

Magnetic Structure Determination and Refinement

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Content:

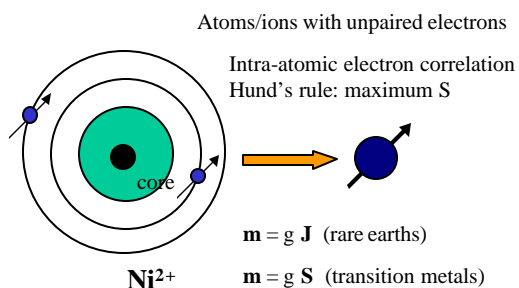
- What's and why magnetic structures
- Formalism to describe magnetic structures
- Magnetic neutron scattering
- Magnetic structure determination:
 - Indexing: *SuperCell*
 - Symmetry Analysis: *BasIreps*
 - Simulated Annealing: *FullProf*
- Examples: practical sessions

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Ions with intrinsic magnetic moments



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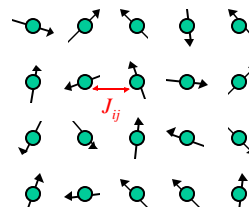
What is a magnetic structure? (1)

Paramagnetic state:

Snapshot of magnetic moment configuration

$$E_{ij} = -J_{ij} \mathbf{S}_i \cdot \mathbf{S}_j$$

$$\langle \mathbf{S}_i \rangle = 0$$



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What is a magnetic structure? (2)

Ordered state: Anti-ferromagnetic

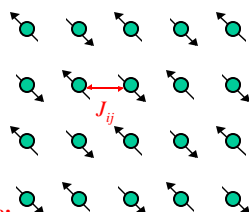
Small fluctuations (spin waves) of the configuration

$$E_{ij} = -J_{ij} \mathbf{S}_i \cdot \mathbf{S}_j$$

$$\langle \mathbf{S}_i \rangle \neq 0$$

Magnetic structure:

Quasi-static configuration of magnetic moments

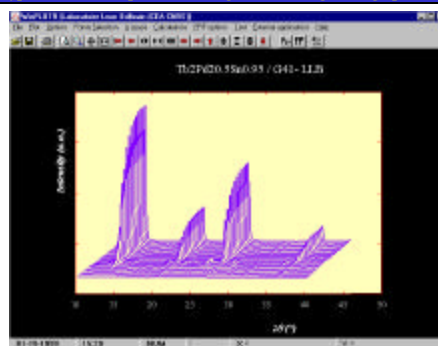


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Typical diffraction patterns of magnetic ordering setting up (G4.1)



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Formalism to describe magnetic structures

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Magnetic structures Magnetic moment of each atom: Fourier series

$$\mathbf{m}_{lj} = \sum_{\{\mathbf{k}\}} \mathbf{S}_{\mathbf{k}j} \exp\{-2\pi i \mathbf{k} \mathbf{R}_l\}$$

Necessary condition for real $\mathbf{m}_{lj} \Rightarrow \mathbf{S}_{-\mathbf{k}j} = \mathbf{S}_{\mathbf{k}j}^*$

$$\mathbf{R}_{lj} = \mathbf{R}_l + \mathbf{r}_j = l_1 \mathbf{a} + l_2 \mathbf{b} + l_3 \mathbf{c} + x_j \mathbf{a} + y_j \mathbf{b} + z_j \mathbf{c}$$

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Examples of Fourier coefficients for simple magnetic structures

Single propagation vector
 $\mathbf{k} = (0,0,0)$ or $\mathbf{k} = 1/2 \mathbf{H}$

$$\mathbf{m}_{lj} = \sum_{\{\mathbf{k}\}} \mathbf{S}_{\mathbf{k}j} \exp\{-2\pi i \mathbf{k} \mathbf{R}_l\} = \mathbf{S}_{\mathbf{k}j} (-1)^{n(l)}$$

REAL Fourier coefficients \equiv magnetic moments

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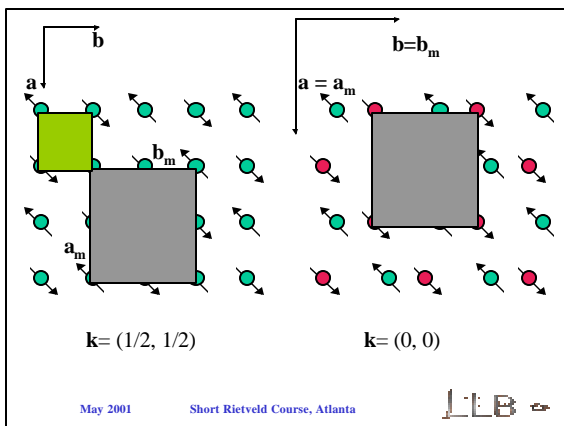
The simplest case Single propagation vector $\mathbf{k} = (0,0,0)$

$$\mathbf{m}_{lj} = \sum_{\{\mathbf{k}\}} \mathbf{S}_{\mathbf{k}j} \exp\{-2\pi i \mathbf{k} \mathbf{R}_l\} = \mathbf{S}_{\mathbf{k}j}$$

- The magnetic structure may be described within the crystallographic unit cell
- Magnetic symmetry similar to conventional crystallography plus time reversal operator.

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Fourier coefficients of sinusoidal structures \mathbf{k} interior of the Brillouin zone (pair \mathbf{k} , $-\mathbf{k}$) Real $\mathbf{S}_{\mathbf{k}}$, or imaginary component in the same direction as the real one

$$\mathbf{m}_{lj} = \mathbf{S}_{\mathbf{k}j} \exp i\{-\mathbf{k} \mathbf{R}_l\} + \mathbf{S}_{-\mathbf{k}j} \exp i\{+\mathbf{k} \mathbf{R}_l\}$$

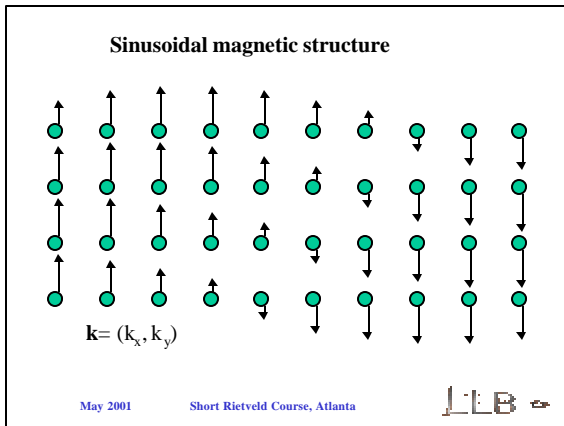
$$\mathbf{S}_{\mathbf{k}j} = \frac{1}{m_j} \sum_j \mathbf{u}_j \exp i\{-2\pi i \mathbf{f}_{kj}\}$$

$$\mathbf{m}_{lj} = m_j \mathbf{u}_j \cos\{2\pi(\mathbf{k} \mathbf{R}_l + \mathbf{f}_{kj})\}$$

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Fourier coefficients of helical structures
 \mathbf{k} interior of the Brillouin zone
Real component of $\mathbf{S}_{\mathbf{k}}$ perpendicular to the imaginary component

$$\mathbf{S}_{\mathbf{k}j} = \frac{1}{2} [m_{uj} \mathbf{u}_j + i m_{vj} \mathbf{v}_j] \exp\{-2p i f_{kj}\}$$

$$\mathbf{m}_{ij} = m_{uj} \mathbf{u}_j \cos\{2p(\mathbf{kR}_i + \mathbf{f}_{ij})\} + m_{vj} \mathbf{v}_j \sin\{2p(\mathbf{kR}_i + \mathbf{f}_{ij})\}$$

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Magnetic neutron scattering

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Magnetic scattering of neutrons

Dipolar interaction (\mathbf{m}_i, \mathbf{m}): vector scattering amplitude

$$\mathbf{a}_M(\mathbf{Q}) = \frac{1}{2} r g f(\mathbf{Q}) \left\{ \mathbf{m} - \frac{\mathbf{Q} (\mathbf{m} \cdot \mathbf{Q})}{Q^2} \right\}$$

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Magnetic scattering of neutrons

$$\mathbf{a}_M(\mathbf{Q}) = \frac{1}{2} r g f(\mathbf{Q}) \left\{ \mathbf{m} - \frac{\mathbf{Q} (\mathbf{m} \cdot \mathbf{Q})}{Q^2} \right\} = p f(\mathbf{Q}) \mathbf{m}_{\perp}$$

$p = 0.2696 \cdot 10^{-12} \text{ cm}$

$$f(\mathbf{Q}) = \int \mathbf{r}_m(\mathbf{r}) \exp(i \mathbf{Q} \cdot \mathbf{r}) d^3 \mathbf{r}$$

Only the perpendicular component of \mathbf{m} to $\mathbf{Q} = 2\pi\mathbf{h}$ contributes to scattering

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Magnetic scattering:

Intensity (non-polarised neutrons)

$$I_h = N_h N_h^* + \mathbf{M}_{\perp h} \cdot \mathbf{M}_{\perp h}^*$$

Magnetic interaction vector:

\mathbf{U} Scattering vector

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Magnetic structures

$$\mathbf{m}_{lj} = \sum_{\{\mathbf{k}\}} \mathbf{S}_{\mathbf{k}j} \exp\{-2\mathbf{p} \cdot i\mathbf{k}\mathbf{R}_l\}$$

The magnetic structure factor:

$$\mathbf{M}(\mathbf{h}) = p \sum_{j=1}^n O_j f_j(\mathbf{h}) T_j \sum_s M_{js} \mathbf{S}_{\mathbf{k}j} \exp\{2\mathbf{p}i[(\mathbf{H}+\mathbf{k})\{\mathbf{t}\}_s \mathbf{r}_j - \mathbf{y}_{\mathbf{k}js}]\}$$

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$$\mathbf{y}_{\mathbf{k}js} = \Phi_{\mathbf{k}j} + \mathbf{f}_{\mathbf{k}js}$$

$\Phi_{\mathbf{k}j}$ is a phase factor which is not determined by symmetry.

$\mathbf{f}_{\mathbf{k}js}$ is a phase factor determined by symmetry.

$$\mathbf{S}_{\mathbf{k}js} = M_{js} \mathbf{S}_{\mathbf{k}j} \exp\{-2\mathbf{p}i\mathbf{f}_{\mathbf{k}js}\}$$

The matrices M_{js} and phases $\mathbf{f}_{\mathbf{k}js}$ can be deduced from the atomic basis functions.

$$\mathbf{S}_{\mathbf{k}js} = \sum_{nl} C_{nl}^n \mathbf{S}_{nl}^{\mathbf{k}n}(js)$$

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The magnetic structure factor:

$$\mathbf{M}(\mathbf{h}) = p \sum_{j=1}^n O_j f_j(\mathbf{h}) T_j \sum_s \mathbf{S}_{\mathbf{k}js} \exp\{2\mathbf{p}i[(\mathbf{H}+\mathbf{k})\{\mathbf{t}\}_s \mathbf{r}_j - \Phi_{\mathbf{k}j}]\}$$



$$\mathbf{M}(\mathbf{h}) = p \sum_{j=1}^n O_j f_j(\mathbf{h}) T_j \sum_{nl} C_{nl}^n \sum_s \mathbf{S}_{nl}^{\mathbf{k}n}(js) \exp\{2\mathbf{p}i[\mathbf{h}_s \mathbf{r}_j - \Phi_{\mathbf{k}j}]\}$$

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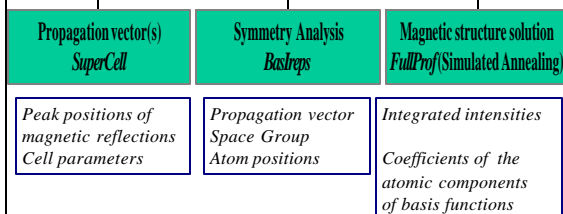
Magnetic structure determination

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Steps for Magnetic structure determination from Neutron Powder Diffraction



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The Program *SuperCell*

(distributed within *WinPLOTR*)

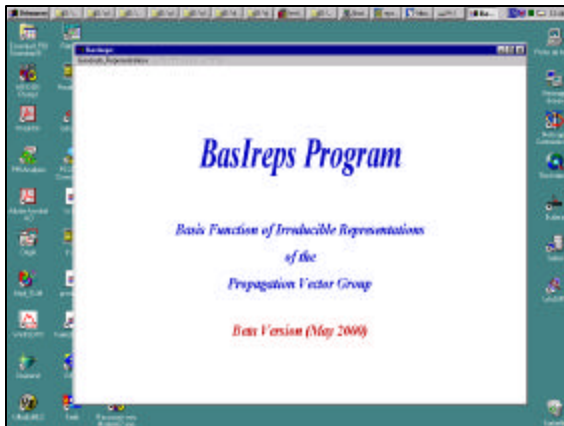
Program: *SuperCell* (J.Rodríguez-Carvajal, LLB-December-1998)

- This program can be used to index superstructure reflections from a powder diffraction pattern.
- The first approach consist in **searching the best "magnetic unit cell"** compatible with a set of observed SUPERSTRUCTURE lines in the powder diffraction pattern.
- If the **first approach fails** to give a suitable solution, the superstructure may be incommensurate and a **direct search for the propagation vector and one of its harmonics have to be used.**

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BasIreps provides the basis functions (normal modes) of the irreducible representations of the wave-vector group G_k

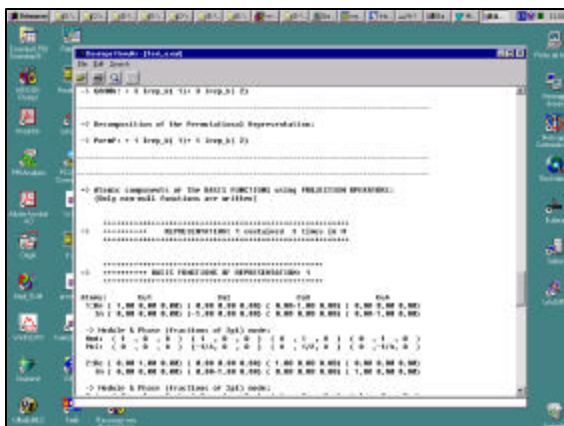
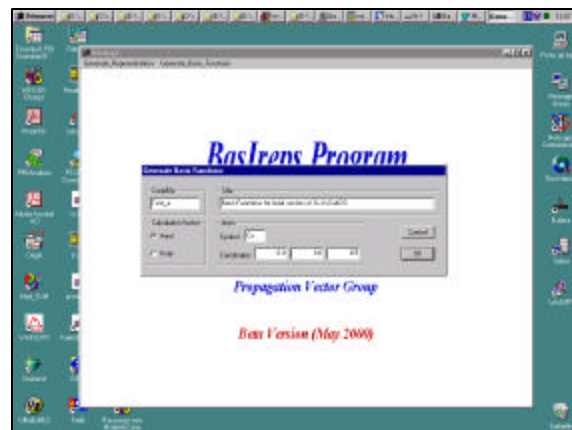
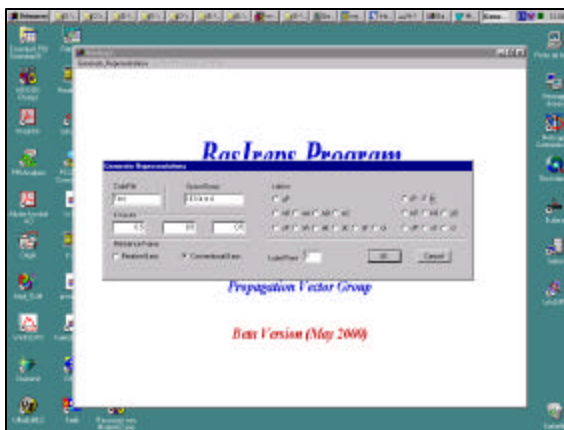
$$\mathbf{m}_{ljs} = \sum_{\{\mathbf{k}\}} S_{\mathbf{k}js} \exp\{-2\mathbf{p} \cdot i\mathbf{k}\mathbf{R}_l\}$$

$$S_{\mathbf{k}js} = \sum_{nI} C_{nI}^n S_{nI}^{\mathbf{k}n}(js)$$

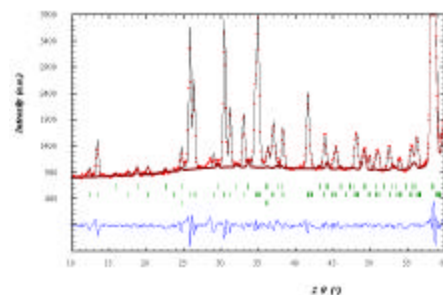
Basis Functions (constant vectors): $S_{nI}^{\mathbf{k}n}(js)$

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Extracting Magnetic Intensities



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Direct space methods:

- Look directly for coefficients of the expansion:

$$S_{kjs} = \sum_{nl} C_{nl}^n S_{nl}^{kn}(js)$$

or components of S_k and phases, explaining the experimental data

- Minimise a reliability factor with respect to the "configuration vector"

$$\mathbf{v} = [C_1, C_2, C_3, C_4, C_5, \dots, C_m]$$

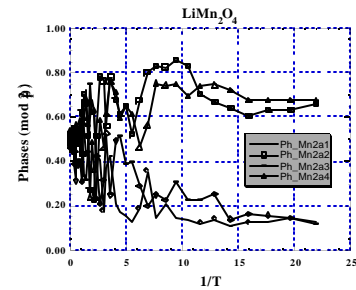
$$R_m(\mathbf{v}) = c \sum_{r=1}^N \left| G_{obs}^2(\mathbf{h}_r) - G_{calc}^2(\mathbf{h}_r, \mathbf{v}) \right|$$

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Behavior of parameters in Simulated Annealing runs

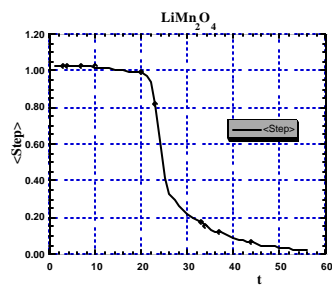


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Average step ...



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